Towards tensor renormalization group study of lattice QCD

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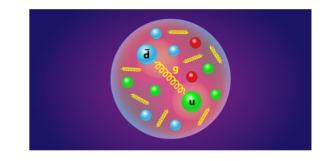






Lattice QCD

- non-perturbative formulation for quantum chromodynamics
- 4D Euclidean SU(3) gauge theory with Nf=2,3
 - Higher dimensions, Non-abelian, Multiple flavors
- MC computation suffers from the sign problem at finite θ and finite density

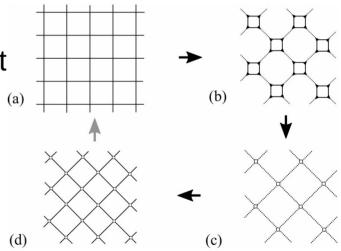


MC computation also suffers from the topology freezing problem toward continuum limit

So far, we do not have a universal method that avoids all the problems...

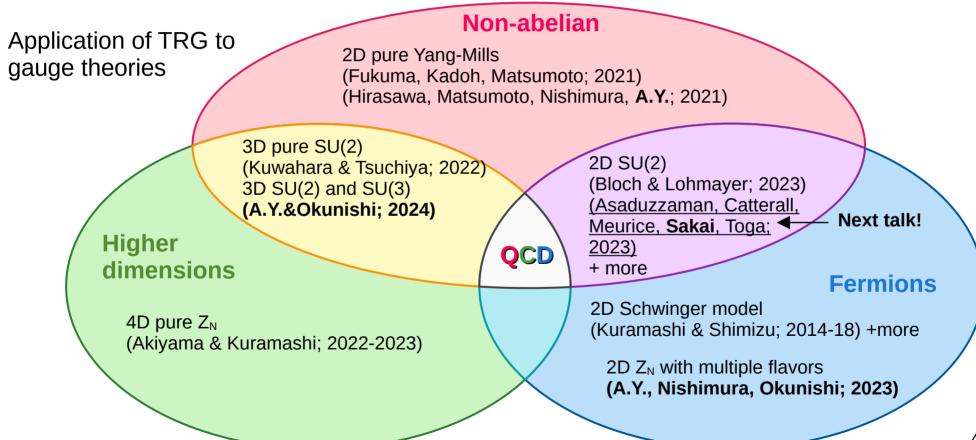
Tensor renormalization group (TRG)

- An alternative to Monte Carlo methods based on coarse graining
- No sampling = No sign problem
- Can access large volumes with log cost
- Can handle fermion/Grassmann numbers directly; Grassmann TRG



[Figures from Okunishi-Nishino-Ueda; 2022]

Progress toward TRG study of Lattice QCD



Challenges

- TRG can be challenging when the local Hilbert space is large
- By that, I meant QCD
 - \rightarrow Multiple fermion flavors ==> dimension \sim exp(k N_f)
 - Non-abelian gauge symmetry ==> Redundancy in the TN

I will talk about my works on these two directions.

Outline

- Part I: Multi-flavor gauge theory
 - The multi-layer formulation
 - Initial tensor compression
 - Result: finite density N_f flavor 2D Z_N theory
- Part II: Armillary sphere formulation
 - Degeneracy in the tensor network
 - Reduced tensor network formulation
 - Result: 3D SU(2) & SU(3) thoery

Part I: Multi-layer construction for multi-flavor gauge theory

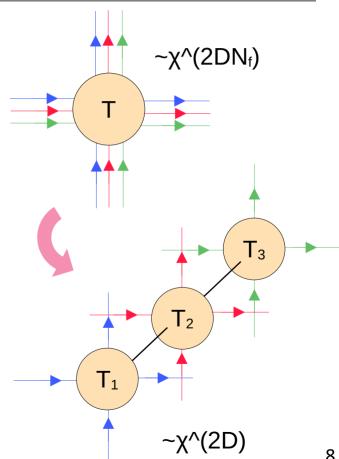
Based on [JHEP11(2023)187], with Jun Nishimura (KEK) and Kouichi Okunishi (Niigata U)

Multi-flavor gauge theory

- The number of tensor legs increases with the number of flavors
- This makes it difficult to consider multi-flavor theory in the tensor network
- This can be avoided by separating flavor d.o.f. from each other

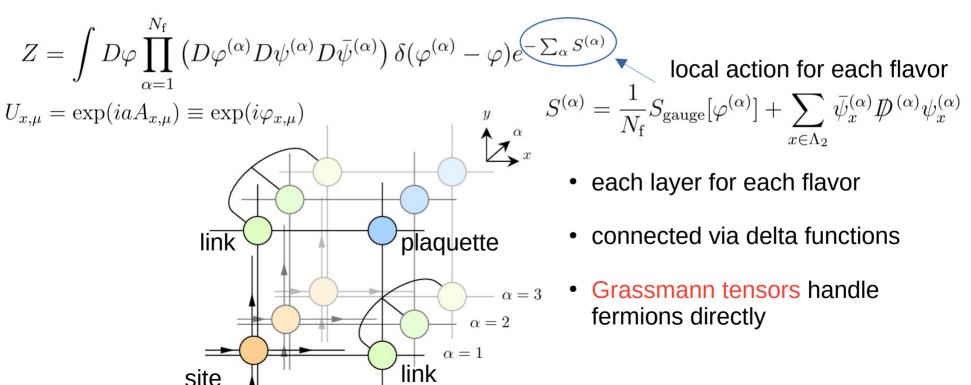
Similar ideas:

- Domain wall fermions (flavor = extra dim)
- MPO-like decomposition (Akiyama; 2023)



Multi-flavor gauge theory

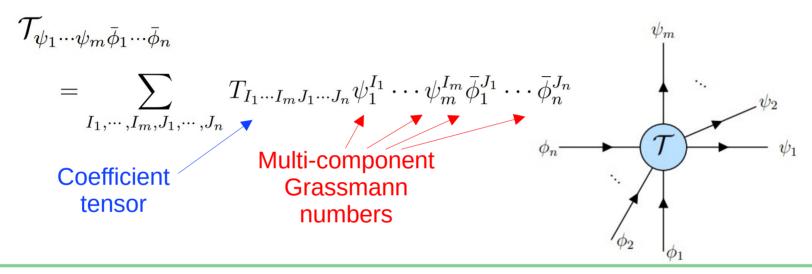
Separate the layers analytically

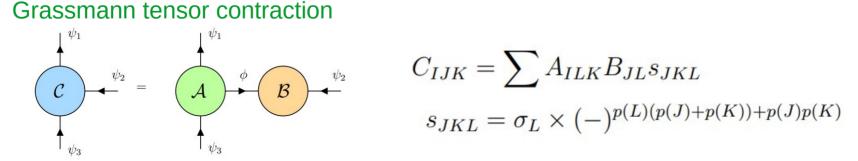


$$S^{(\alpha)} = \frac{1}{N_{\rm f}} S_{\rm gauge}[\varphi^{(\alpha)}] + \sum_{x \in \Lambda_2} \bar{\psi}_x^{(\alpha)} \mathcal{D}^{(\alpha)} \psi_x^{(\alpha)}$$

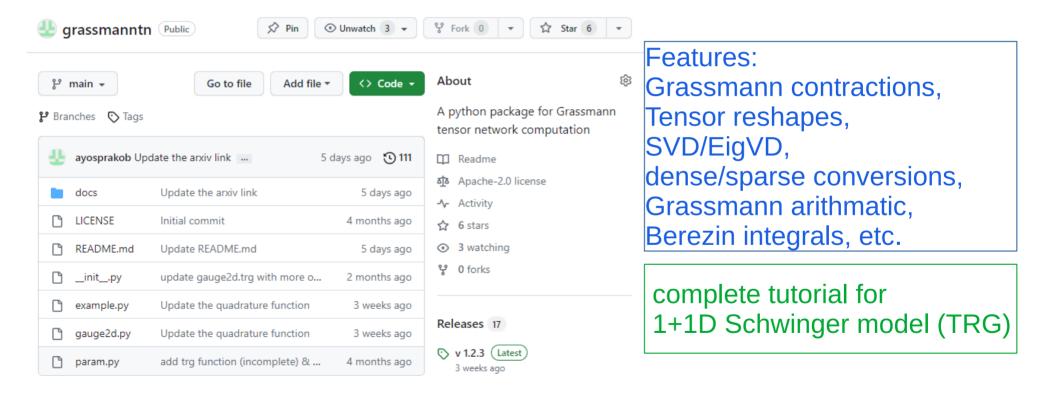
- · each layer for each flavor
- connected via delta functions
- Grassmann tensors handle fermions directly

Quick intro: Grassmann tensors

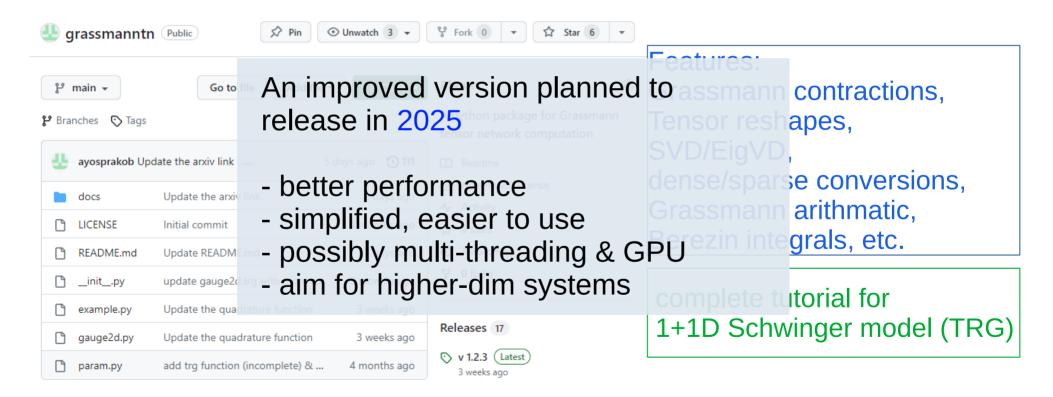




GrassmannTN: a python package for Grassmann TRG / DMRG

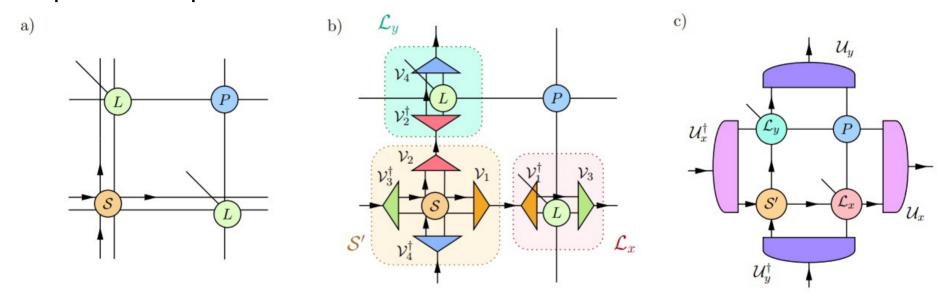


GrassmannTN: a python package for Grassmann TRG / DMRG



Tensor compression

Proposed compression scheme:



Isometries are first applied around the Grassmann tensor S: $a \rightarrow b$ Then another set is applied around the whole tensor: $b \rightarrow c$

Compression performance

Physical parameters

$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	β	$\tilde{\mu}$	$N_{ m f}$	K	original size	compressed size	compression ratio	D_x	D_y
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	0.0	0.0	1	2	67108864	1024	1.53×10^{-5}	4	_
	0.0	0.0	1	3	3869835264	2304		4	4
$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	0.0	0.0	1	4	68719476736	4096	5.96×10^{-8}	4	4
$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	0.0	0.0	1	5	640000000000	6400	1.00×10^{-9}	4	4
$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	2.0	0.0	1	2	67108864	16384	2.44×10^{-4}	8	8
	2.0	0.0	2	2	67108864	16384	2.44×10^{-4}	8	8
$\begin{bmatrix} 2.0 & 3.0 & 2 & 2 & 67108864 & 16384 & 2.44 \times 10^{-4} & 8 & 8 \end{bmatrix}$	2.0	3.0	1	2	67108864	16384	2.44×10^{-4}	8	8
2.0 0.0 2 2 0.100001	2.0	3.0	2	2	67108864	16384	2.44×10^{-4}	8	8

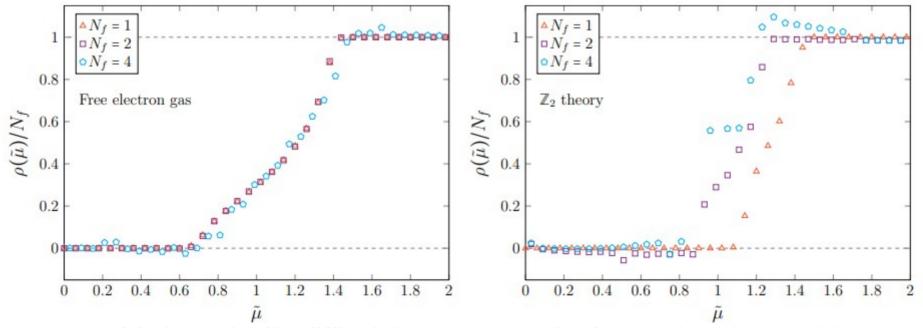
All of these compressions are done without any spectrum truncation! Such high performance is due to the sparse nature of fermions.

Finite density and Silver Blaze

Flavors: HOTRG (Dcut=64, 32 for Nf=2, 4)

Space-time: Levin-Nave TRG (Dcut=64)

Silver Blaze is reproduced up to $N_f = 4$



This is typically difficult in Monte Carlo due to the sign problem!

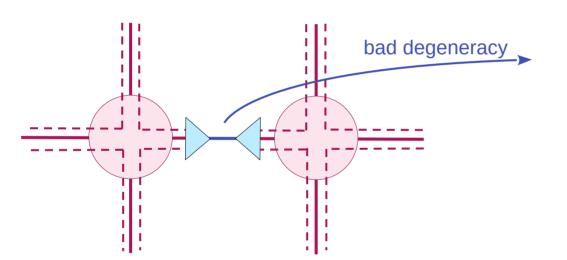
Part II: Armillary sphere Non-abelian gauge theory in higher dimensions

Based on [PTEP 2024 (2024) 7, 073B05] (Formulation) and [arXiv:2406.16763] (Numerical) with **Kouichi Okunishi** (Niigata U)

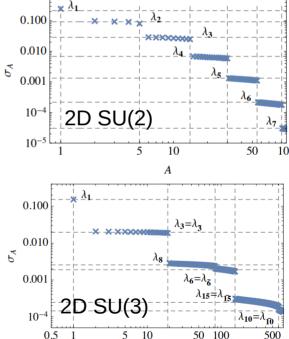
Why is non-abelian tensor network difficult?

Internal symmetry in SU(N) is a redundancy in the tensor network

that cannot be truncate by an SVD



The entanglement structure is nonlocal...



Figures from [Fukuma-Kadoh-Matsumoto; 2021]

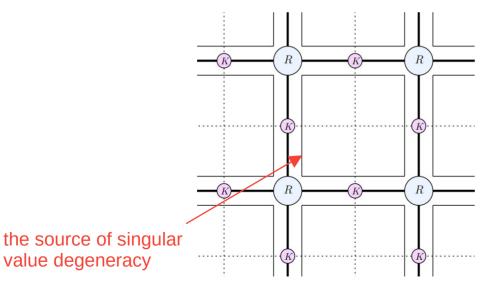
Character expansion

value degeneracy

 Lesson from 1+1D: the (matrix) index loops can be traced out if we use character expansion

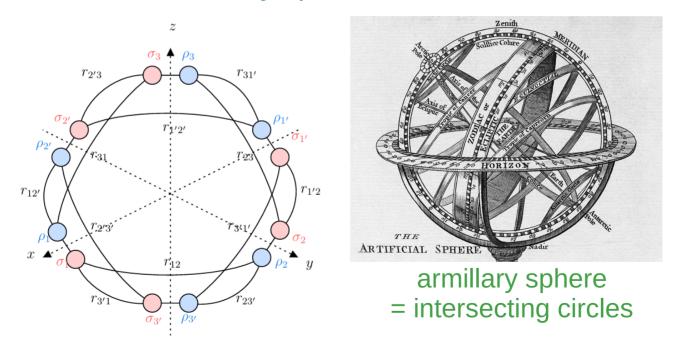
[Hirasawa, Matsumoto, Nishimura, A.Y.; 2021]

Degeneracy is completely eliminated



Can we do the same thing for any dimension?

Yes! There is a similar closed network in any dimension Which we call the armillary sphere



This was first noticed by [Oeckl & Pfeiffer;2001] in the context of the spin foam model.

Step 1: perform character expansion on the Boltzmann weight

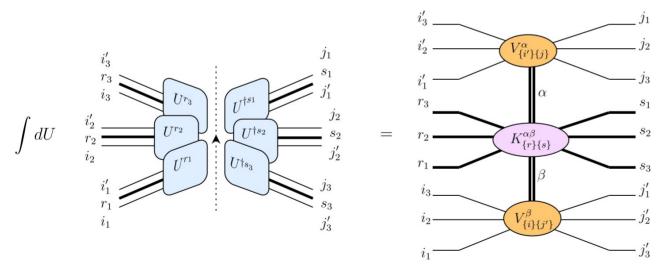
$$e^{\beta\Re\operatorname{tr} P_{n,\mu\nu}} = \sum_{r} f_r \operatorname{tr}_r(U_{n,\mu} U_{n+\hat{\mu},\nu} U_{n+\hat{\nu},\mu}^{\dagger} U_{n,\nu}^{\dagger})$$

$$= \sum_{r} f_r \sum_{i,j,k,l} (U_{n,\mu})_{ij}^r (U_{n+\hat{\mu},\nu})_{jk}^r (U_{n+\hat{\nu},\mu}^{\dagger})_{kl}^r (U_{n,\nu}^{\dagger})_{li}^r$$

$$\operatorname{tr}_r P_{n,\mu\nu} = \underbrace{ (U_{n,\nu}^{\dagger})_{\ell i}^r }_{i} \underbrace{ (U_{n+\hat{\mu},\nu})_{jk}^r (U_{n+\hat{\mu},\nu})_{jk}^r }_{i} \underbrace{ (U_{n+\hat{\mu},\nu})_{jk}^r }_{j} \underbrace{ (U_{n+\hat{\mu},\nu})_{jk}^r }_{i} \underbrace{ (U_{n+\hat{\mu},\nu})_{jk}^r }_{j} \underbrace{ (U_{n+\hat{\mu},\nu})_{jk}^r }_{i} \underbrace{ (U_{n+\hat{\mu},\nu})_{jk}^r }_{i}$$

Step 2: perform group integral on each link variable

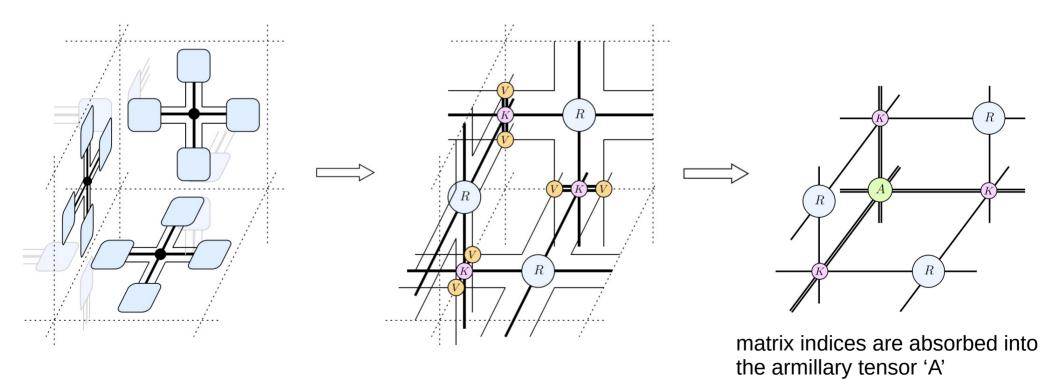
$$\int dU_{n,\mu}(U_{n,\mu})_{i_1i'_1}^{r_1}\cdots(U_{n,\mu})_{i_{d-1}i'_{d-1}}^{r_{d-1}}(U_{n,\mu}^{\dagger})_{j_1j'_1}^{s_1}\cdots(U_{n,\mu}^{\dagger})_{j_{d-1}j'_{d-1}}^{s_{d-1}} = \sum_{\hat{r}\in D_{\otimes r}}\sum_{\hat{s}\in D_{\otimes s}}\sum_{\hat{\iota},\hat{\jmath}}\frac{1}{\dim(\hat{r})}C_{\{r\}\{i\}}^{\hat{r}\hat{\iota}}C_{\{s\}\{j\}}^{\hat{s}\hat{\jmath}}C_{\{s\}\{j'\}}^{\hat{s}\hat{\iota}}\delta(\hat{r}\simeq\hat{s})$$



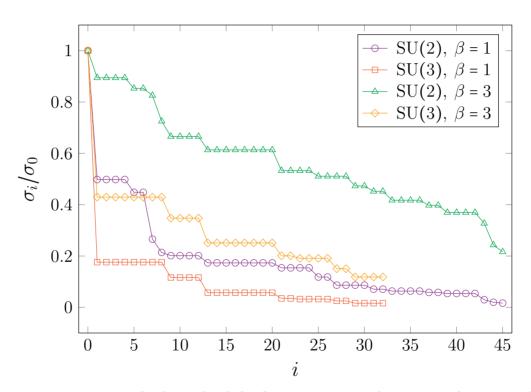
$$\begin{split} V^{\alpha}_{\{i\}\{j'\}} &= \delta(\hat{r} \simeq \hat{s}) \sum_{\hat{i}} C^{\hat{r}\hat{i}}_{\{r\}\{i\}} C^{\hat{s}\hat{i}}_{\{s\}\{j'\}}, \\ K^{\alpha\beta}_{\{r\}\{s\}} &= \frac{1}{\dim(\hat{r})} \delta_{\hat{r}\hat{r}'} \delta_{\hat{s}\hat{s}'} \delta(\hat{r} \in D_{\otimes r}) \delta(\hat{s} \in D_{\otimes s}); \end{split}$$

Note: matrix indices (thin lines) are neatly separated into two layers

Step 3: Contract the matrix indices



Result: singular value spectrum



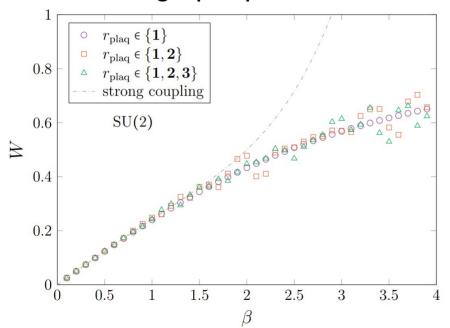
Singular value spectrum of the initial tensor do not have large degeneracy

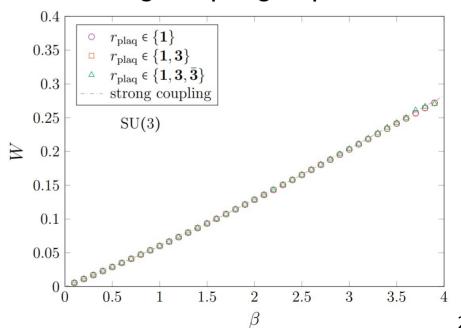
Result: average plaquette @ zero temperature

pure 2+1D SU(2) and SU(3) gauge theory

ATRG; $V = 16^3$; $D_{cut} = 16$

Average plaquette – consistent with strong coupling expansion

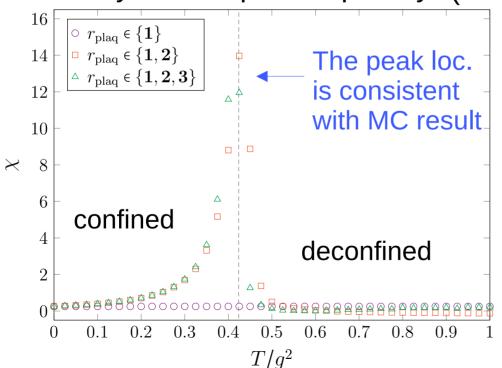


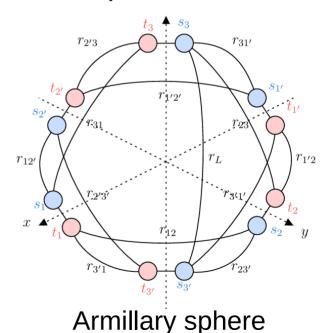


Result: deconfinement @ finite temperature

TRG; $V = 1 \times 1024^2$; $D_{cut} = 64$

Polyakov loop susceptibility (with induced ssb)





with polyakov loop

Summary

- TRG is a promising methods for studying lattice theories
- We address 2 challenging aspects toward lattice QCD
 - Multiple fermion flavors can be handled with Grassmann Tensors with multi-layer construction
 - Degeneracy in non-abelian tensor network can be eliminated with the armillary sphere technique

Future prospects

- Can we reduce the tensor network without character expansion?
 (Some variation of Gilt-TNR?) [Hauru, Delcamp, Mizera; 2017]
- Armillary sphere method with matter fields?
- More in-depth analysis (physical interpretation?)
- 4D gauge theory + theta term
- Etc.