

Quantum algorithms for simulating the Schwinger model

Based on Quantum 8, 1474 (2024), arXiv:2311.17388

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Today's Content

1. Quantum algorithms for simulating quantum systems
 - Why do we use quantum computers?
 - Trotter formula (Conventional method)
 - Quantum singular value transformation (Recent method)
2. Application to particle physics
 - Our recent work:
efficient quantum algorithm for simulating the Schwinger model

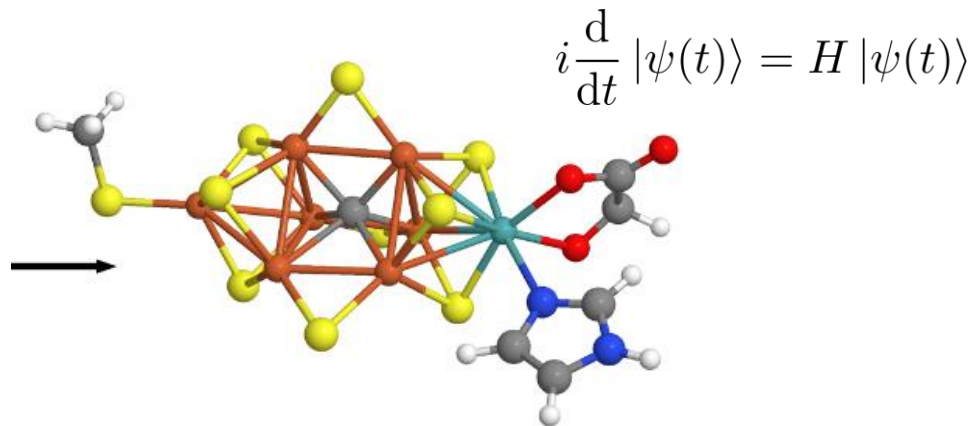
Quantum algorithms for simulating quantum systems

Introduction to quantum computers

Computers which utilize the properties of quantum systems

- Quantum computers are expected to obtain (**exponential**) **quantum speed ups** over classical computers for some problems.

- **Quantum simulation** S. Lloyd, Science, (1996).



M. Reiher, et al. PNAS, (2017).

- Condensed matter physics
- Quantum chemistry
- High energy physics

- **Linear algebraic problems**

A. W. Harrow, Phys. Rev. Lett., (2009).



$$A |x\rangle = |b\rangle$$

- Factoring
- Fluid dynamics simulation
- Machine learning
- Data analysis

Introduction to quantum computers

Components of a quantum computer

Qubit

- 2-dim quantum states.

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

$$|\alpha|^2 + |\beta|^2 = 1$$

- multiple qubit states are described by tensor product.

$$|\psi_1\rangle \otimes |\psi_2\rangle \otimes \cdots \otimes |\psi_n\rangle$$

- n -qubit states are equivalent to 2^n -dim vectors.

$$|\psi\rangle = \sum_{n=0}^{2^n-1} \psi_n |n\rangle = \begin{bmatrix} \psi_0 \\ \psi_1 \\ \vdots \\ \psi_{2^n-1} \end{bmatrix}$$

Unitary gate

- Quantum states are evolved under unitary operations.

- n -qubit unitary gates are $2^n \times 2^n$ unitary matrices.

$$\begin{bmatrix} \psi'_0 \\ \psi'_1 \\ \vdots \\ \psi'_{2^n-1} \end{bmatrix} = \begin{bmatrix} & & & \\ & U & & \\ & & & \\ & & & \end{bmatrix} \begin{bmatrix} \psi_0 \\ \psi_1 \\ \vdots \\ \psi_{2^n-1} \end{bmatrix}$$

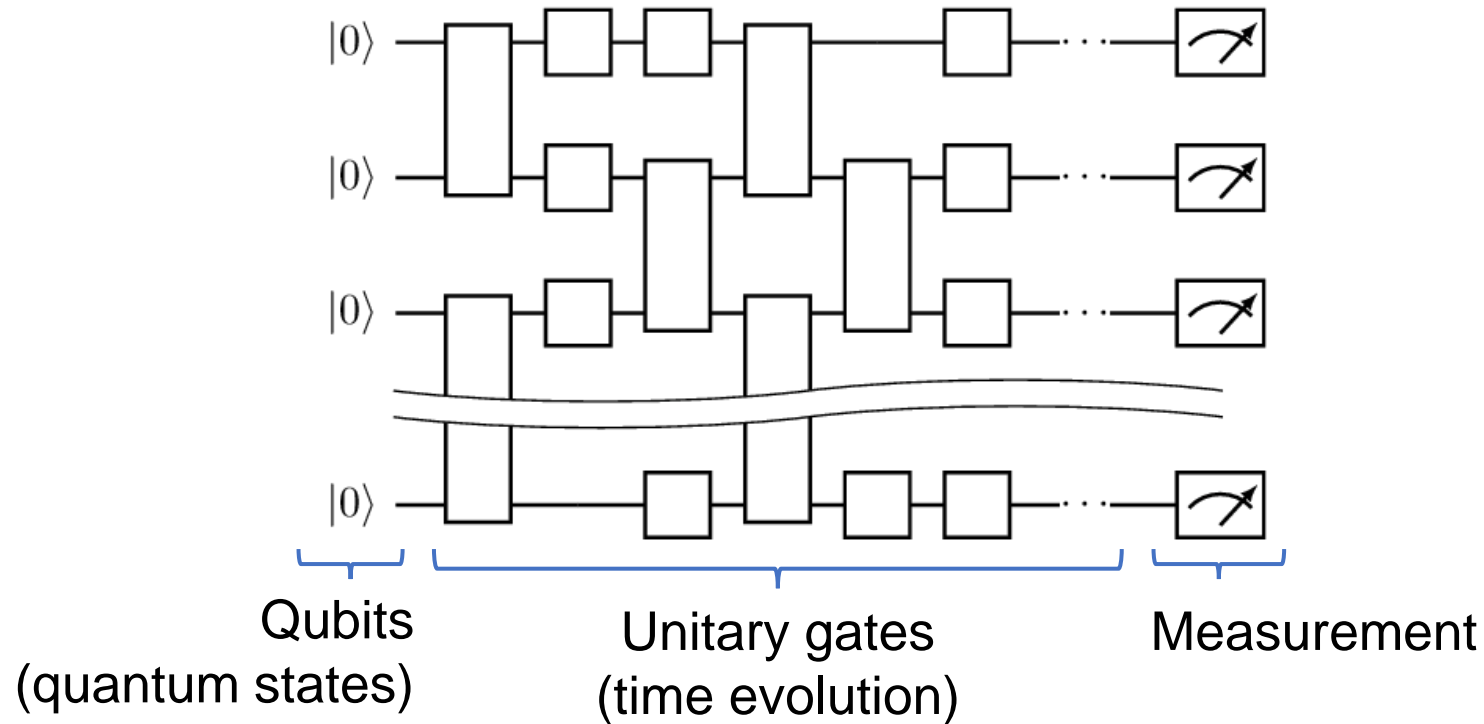
- $\{H, \text{CNOT}, T\}$ gate set is universal and usually used.

Measurement

- We have to measure the evolved state to extract information (solution).
- Measurement protocol performs the task efficiently.
 - ✓ Which qubits?
 - ✓ Which basis?
 - ✓ How many samples?

Introduction to quantum computers

Quantum circuits – diagrams of quantum computing procedures



Quantum algorithms \approx Constructing a quantum circuit to solve a given problem

Hamiltonian simulation

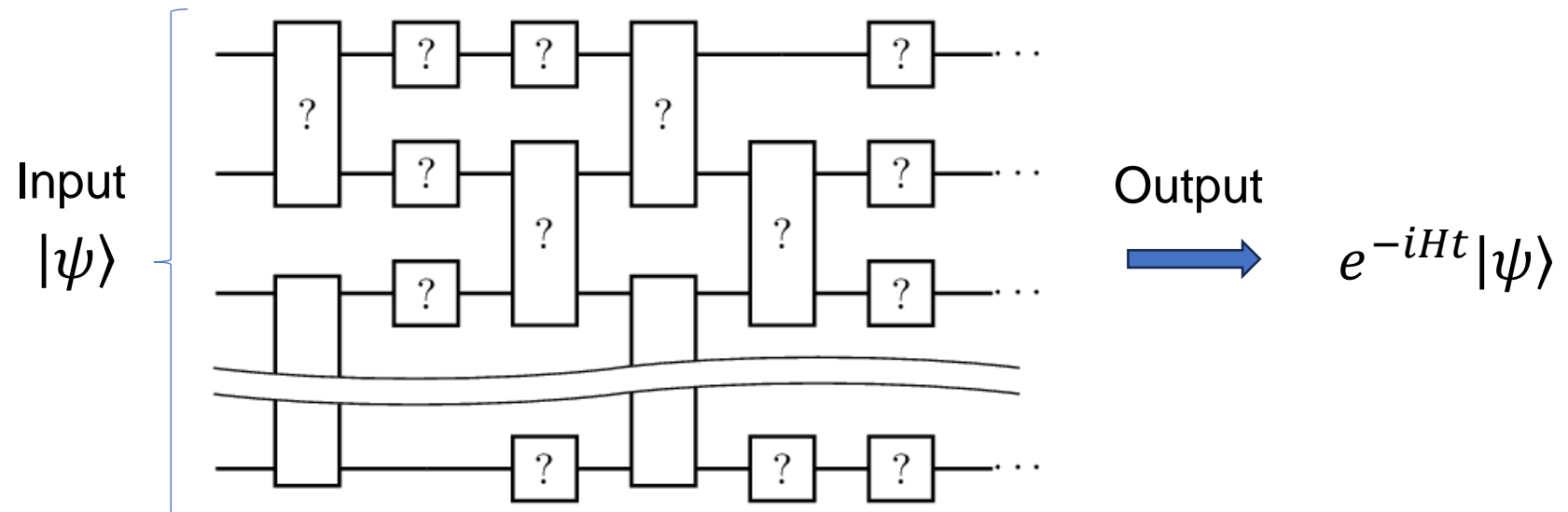
- Hamiltonian simulation is one of the promising task with **exponential quantum advantage**.

S. Lloyd, Science, (1996).

- Problem statement:

Input: a Hamiltonian H and an initial quantum state $|\psi\rangle$.

Output: the evolved quantum state $e^{-iHt}|\psi\rangle$.



How to construct a quantum circuit for simulating a Hamiltonian?

Trotter formula (Conventional method)

S. Lloyd, Science, (1996).

Divide the whole evolution into short time evolution of each term

➤ Given a Hamiltonian $A + B$, we can implement $e^{-i(A+B)t}$ if we have e^{-iAt} and e^{-iBt} .

➤ (first-order) Trotter formula:

$$e^{-i(A+B)t} = \left(e^{-iAt/r} e^{-iBt/r} \right)^r + O\left(\frac{\|[A, B]\|t^2}{r} \right)$$

A. M. Childs, et al. Phys. Rev. X, (2021).

➤ To achieve the total error ε , we set

$$r = O\left(\frac{\|[A, B]\|t^2}{\varepsilon} \right)$$

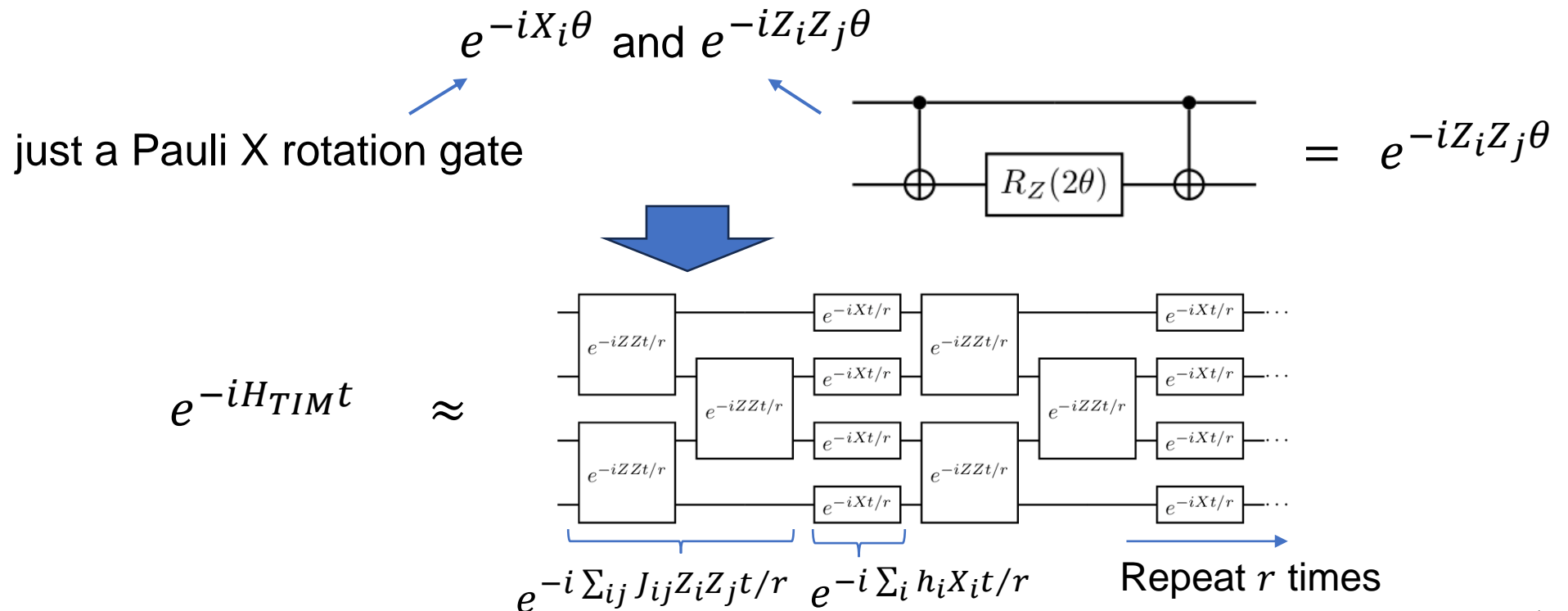
Trotter formula (Conventional method)

S. Lloyd, Science, (1996).

Example: transverse field Ising model

$$H_{TIM} = \sum_{i,j=1}^n J_{ij} Z_i Z_j + \sum_{i=1}^n h_i X_i$$

➤ Trotter formula implies that we only need to implement



Quantum singular value transformation (Recent method)

A. Gilyen, et al, STOC (2019)

- Trotter formula achieves the complexity of $O(\|[A, B]\|t^2/\varepsilon)$.

Can we further improve this complexity?



Yes!

- Quantum singular value transformation (QSVT) achieves $O(\|H\|t + \log(1/\varepsilon))$.
 - Provably optimal complexity in terms of $\|H\|, t, 1/\varepsilon$. A. Gilyen, et al, STOC (2019)
- QSVT implements a polynomial transformation of any matrix.
 - Today we focus on a Hermitian matrix H .

Quantum singular value transformation (Recent method)

A. Gilyen, et al, STOC (2019)

- **Block-encoding:** embed a Hamiltonian H into a unitary operator U .

S. Chakraborty, et al, ICALP, (2019)

$$U = |0^b\rangle\langle 0^b| \otimes H + \dots = \begin{bmatrix} H & \cdot \\ \cdot & \cdot \end{bmatrix}$$

- Let parametrized reflection operator $R(\phi)$ as:

$$R(\phi) = e^{i\phi}|0^b\rangle\langle 0^b| \otimes I + e^{-i\phi}(I - |0^b\rangle\langle 0^b|) \otimes I$$

- This operator can be implemented using multi-controlled gates.

- Surprisingly, a sequence of U and $R(\phi)$ provides a polynomial transformation of H .

$$R(\phi_d)U \cdots R(\phi_2)U^\dagger R(\phi_1)UR(\phi_0) = \begin{bmatrix} P(H) & \cdot \\ \cdot & \cdot \end{bmatrix} \quad d: \text{degree of polynomial } P$$

- QSVT with polynomial approximation of e^{-iHt} provides a new Hamiltonian simulation method.

$$e^{-iHt} \approx \sum_{k=0}^d \beta_k H^k \quad \text{with } d = O(\|H\|t + \log(1/\varepsilon))$$

Block-encoding

Example: linear combination of unitary operators

$$H = \sum_{l=0}^{L-1} \alpha_l U_l$$

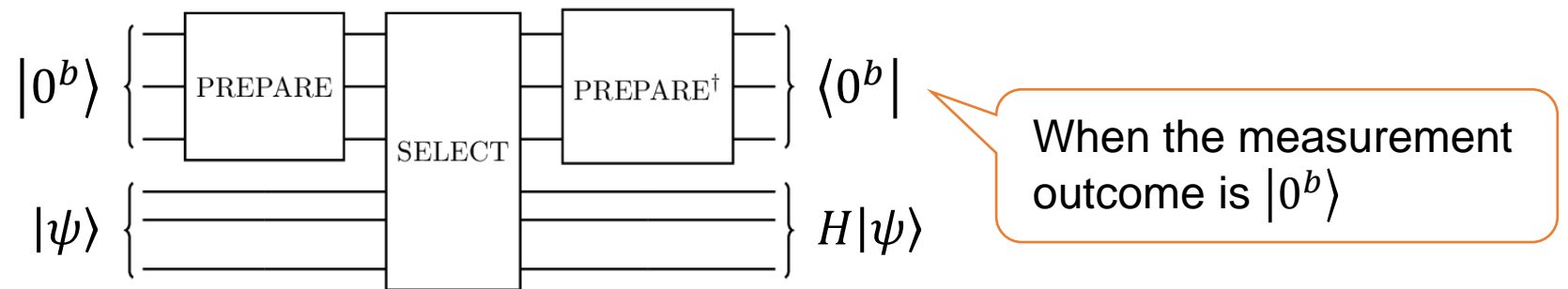
- Assume $a_l > 0$ and $\sum_l a_l = 1$.

➤ Let PREPARE operator and SELECT operator as:

$$\text{PREPARE } |0^b\rangle = \sum_{l=0}^{L-1} \sqrt{\alpha_l} |l\rangle \quad \text{SELECT} = \sum_{l=0}^{L-1} |l\rangle \langle l| \otimes U_l$$

➤ Block-encoding circuit:

R. Babbush, et al. Phys. Rev. X, (2018).



➤ PREPARE and SELECT can be implemented with $O(L)$ gates.

Summary – Hamiltonian simulation

- Trotter formula $e^{-i(A+B)t} \approx (e^{-iAt/r} e^{-iBt/r})^r$
 - The complexity of $O(t^2/\varepsilon)$
 - Simple implementation
- Quantum singular value transformation (QSVT) $e^{-iHt} \approx \sum_{k=0}^d \beta_k H^k$
 - The optimal complexity of $O(t + \log(1/\varepsilon))$
 - Applicable to a block-encoded Hamiltonian

$$R(\phi_d) \overset{\downarrow}{U} \cdots R(\phi_2) \overset{\downarrow}{U}^\dagger R(\phi_1) \overset{\downarrow}{U} R(\phi_0) = \begin{bmatrix} e^{-iHt} & \cdot \\ \cdot & \cdot \end{bmatrix}$$

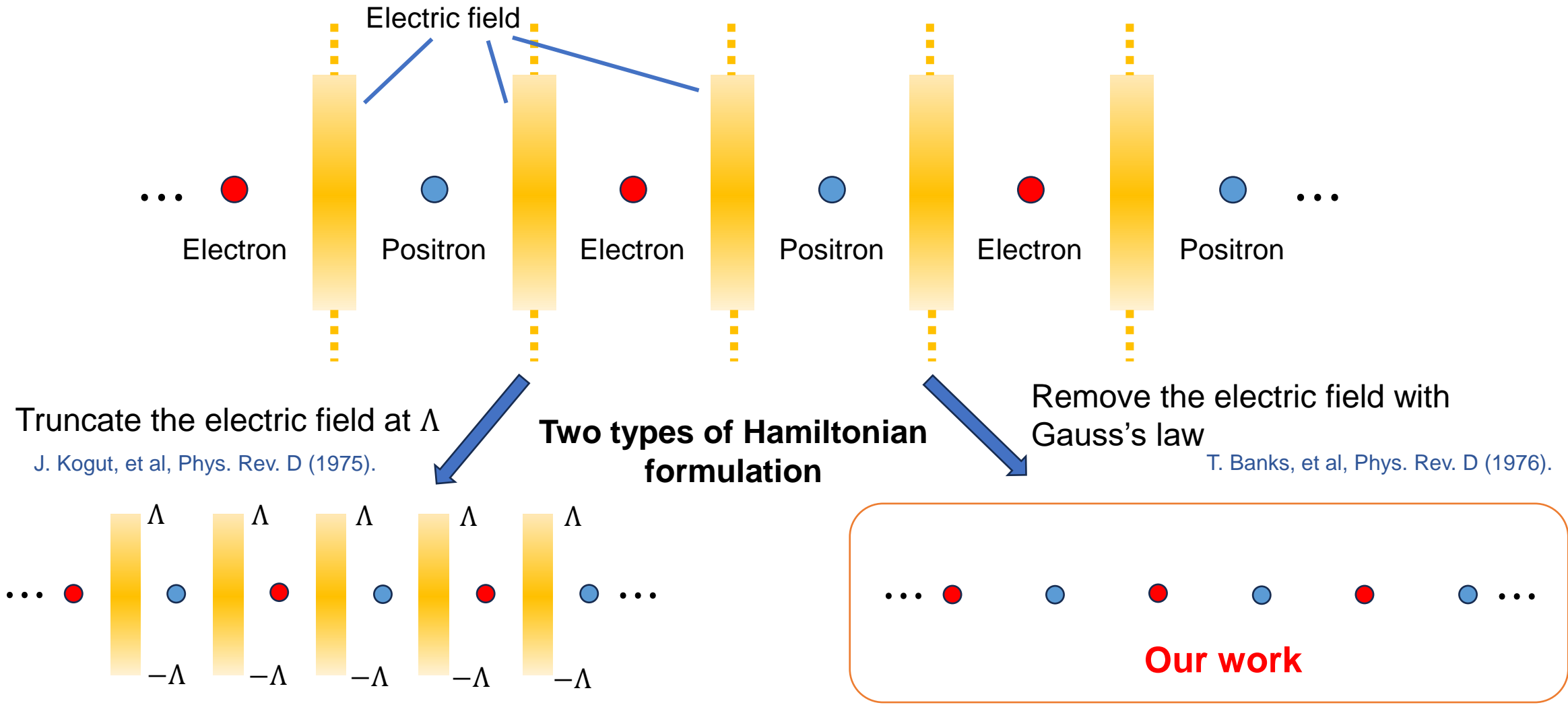
Efficient implementation of block-encoding is important!

Our recent work: efficient quantum algorithm for simulating the Schwinger model

Kazuki Sakamoto, Hayata Morisaki, Junichi Haruna, Etsuko Ito, Keisuke Fujii, Kosuke Mitarai,
“End-to-end complexity for simulating the Schwinger model on quantum computers”,
Quantum 8, 1474 (2024), arXiv:2311.17388

Schwinger model

- one of the simplest yet non-trivial gauge theories



Schwinger model

Previous works (real time evolution e^{-iHt})

- The Hamiltonian formulation without electric field

E. A. Martinez, et al, Nature 534, 516 (2016).

N. H. Nguyen, et al, PRX Quantum 3, 020324 (2022).

- Based on Trotter formula

- System size : N
- Precision : ε
- Evolution time : t

..... $O(N^{4.5}t^{1.5}/\varepsilon^{0.5})$

Our work $\tilde{O}(N^4t + \log(1/\varepsilon))$

- The Hamiltonian formulation which includes electric field

A. F. Shaw, et al, Quantum 4, 306 (2020).

- Based on Trotter formula
- Provides rigorous cost analysis

..... $\tilde{O}(N^{2.5}t^{1.5}/\varepsilon^{0.5})$

Y. Tong, et al, Quantum 6, 816 (2022).

- The smallest query complexity at present
- Needs a huge number of qubits

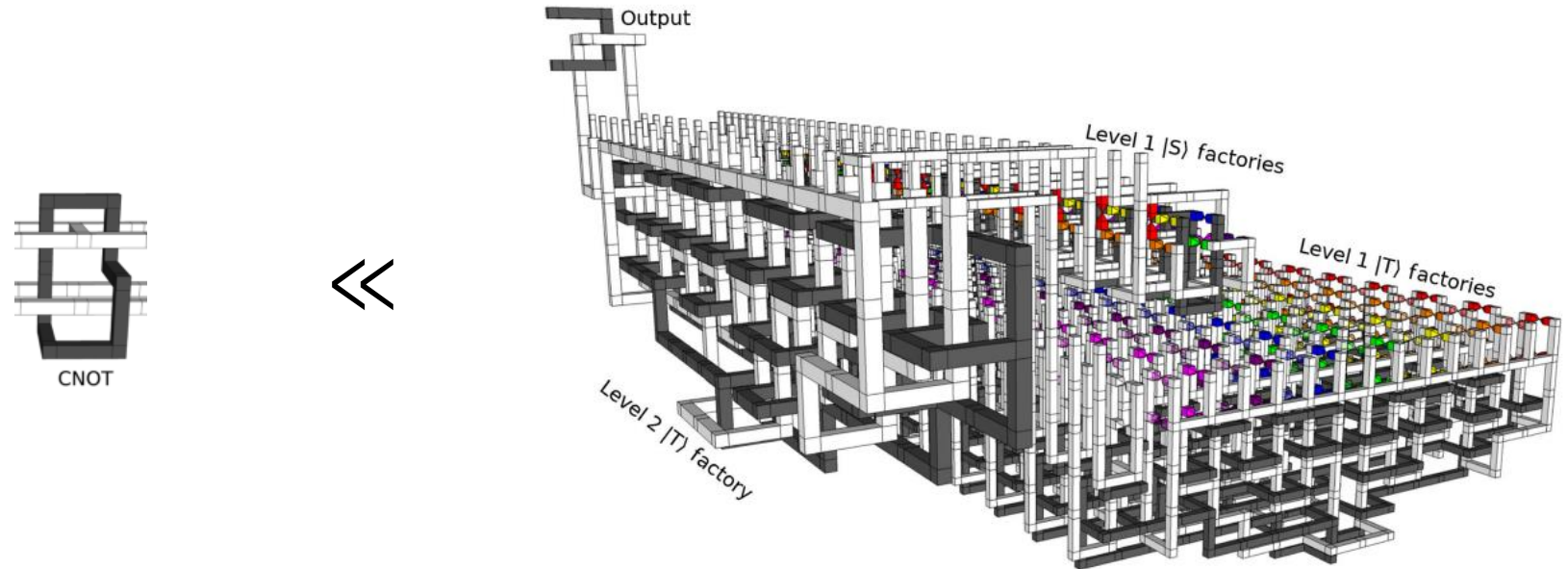
..... $\tilde{O}(Nt \text{ polylog}(1/\varepsilon))$

Rigorous resource estimates without scaling are needed for a comparison.

Runtime is estimated via the number of T gates

T gate is costly in FTQC due to error correction

- An arbitrary unitary operator can be decomposed into Clifford + T gates.
- Clifford gates are easy, but T gates need large space-time overhead.



R. Babbush, et al. Phys. Rev. X, (2018).

The number of T gates dominates actual runtime.

Result – efficient block-encoding

- The Schwinger model Hamiltonian after Jordan-Wigner transformation:

$$H_S = J \sum_{n=0}^{N-2} \left(\sum_{i=0}^n \frac{Z_i + (-1)^i}{2} + \frac{\theta_0}{2\pi} \right)^2 + \frac{w}{2} \sum_{n=0}^{N-2} (X_n X_{n+1} + Y_n Y_{n+1}) + \frac{m}{2} \sum_{n=0}^{N-1} (-1)^n Z_n$$

- Naively, we need $O(N^2)$ T gates for the block-encoding since H_S has $O(N^2)$ terms.

- ◆ **Our implementation requires only $O(N)$ T gates.**

- Decompose the Hamiltonian into several parts as below:

$$H_S = \boxed{\frac{J}{4} \sum_{n=1}^{N-1} \left(\sum_{i=0}^{n-1} Z_i \right)^2} + \boxed{J \frac{\theta}{2\pi} \sum_{n=1}^{N-1} \sum_{i=0}^{n-1} Z_i} + \boxed{J \left(\frac{1}{2} + \frac{\theta}{2\pi} \right) \sum_{n=1}^{N-1} \sum_{i=0}^{n-1} Z_i} + \boxed{\frac{w}{2} \sum_{n=0}^{N-2} X_n X_{n+1}} + \boxed{\frac{w}{2} \sum_{n=0}^{N-2} Y_n Y_{n+1}} + \boxed{\frac{m}{2} \sum_{n=0}^{N-1} (-1)^n Z_n}$$

- Uniform superposition states $\frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |i\rangle$ can be prepared efficiently with $O(\log N)$ T gates.

Y. R. Sanders, et al, PRX Quantum (2020).

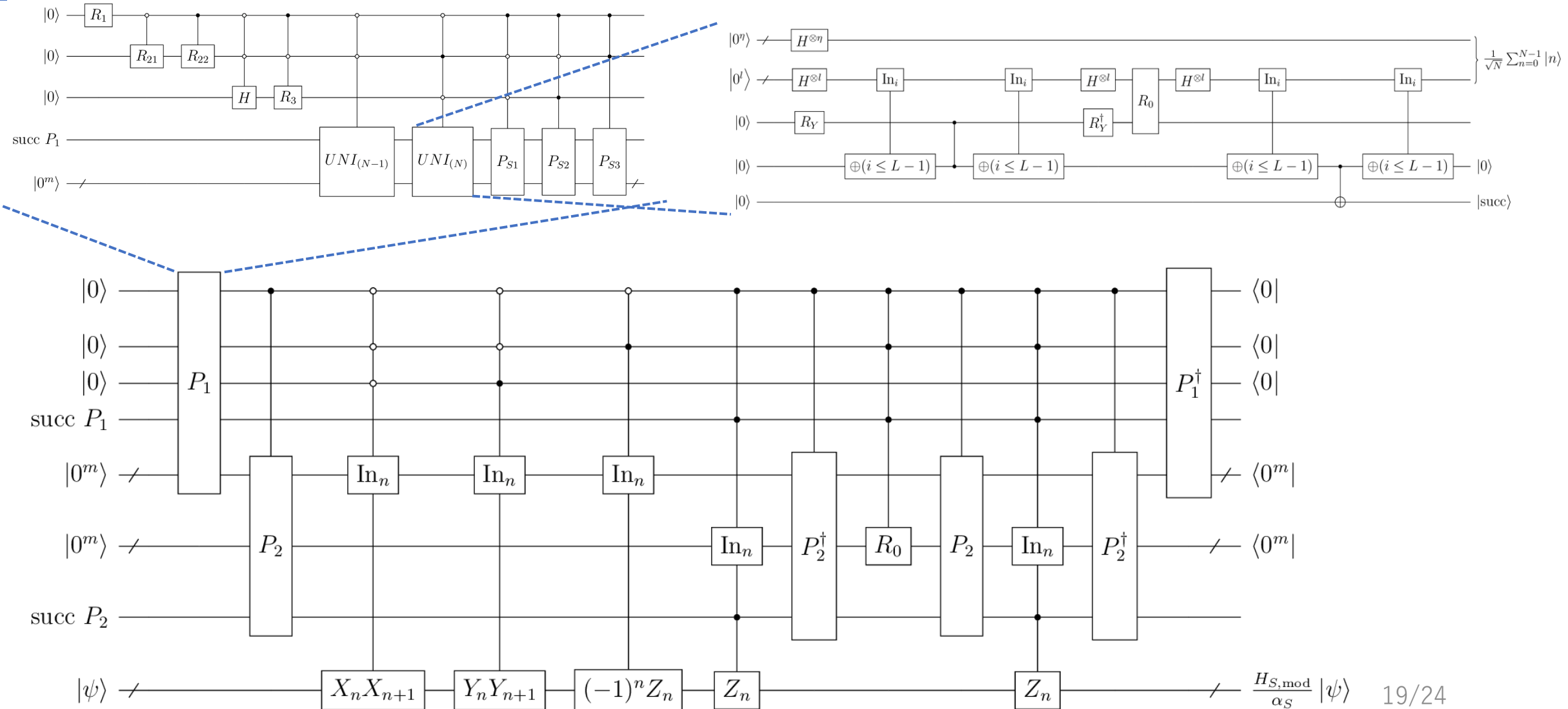
➡ Leading to efficient block-encodings of each term

- Take a linear combination of the block-encodings of each term

A. Gilyen, et al, STOC (2019)

Result – efficient block-encoding

Overall circuit description



Result – resource estimates

Simulating the creation and annihilation of particle pairs

- The ground state of H_S for $J = \theta_0 = w = 0, m = m_0$ is the vacuum state.

$$|\text{vac}\rangle = |1010 \dots\rangle$$

- The state without any particle.

- Evolve $|\text{vac}\rangle$ under H_S for time t and estimate the amplitude of $|\text{vac}\rangle$.

$$\left| \langle \text{vac} | e^{-iH_S t} | \text{vac} \rangle \right| \quad : \text{vacuum persistence amplitude} \quad \text{J. Schwinger, Phys. Rev. (1951).}$$

- Simulation of the creation and annihilation of electron-positron pairs.



**Based on our block-encoding,
how much resource is required to compute this quantity?**

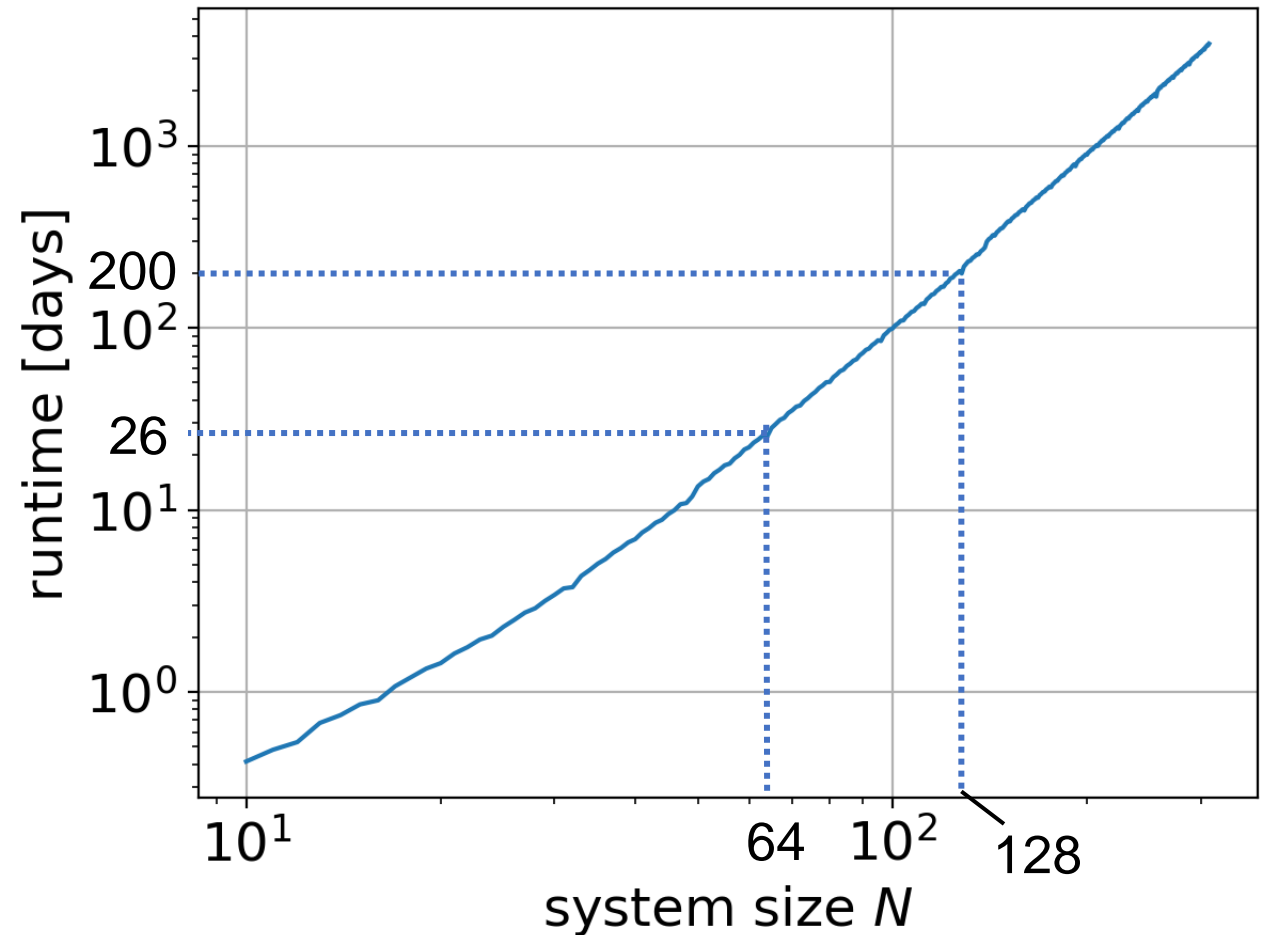
Result – resource estimates

Runtime

- Parameters
 - Precision (additive error) : $\varepsilon = 0.01$
 - Evolution time : $t = 4$
 - T gate consumption rate : 1MHz
 - Lattice spacing : $a = 0.2$
 - electron mass : $m = 0.1$
 - $w = \frac{1}{2a} = 2.5$
 - $J = \frac{g^2 a}{2} = 0.1, (g = 1)$
 - $\theta_0 = \pi$

Examples

System size	Runtime [days]
64	26
128	200



Runtime for calculating the vacuum persistence amplitude.

Result – resource estimates

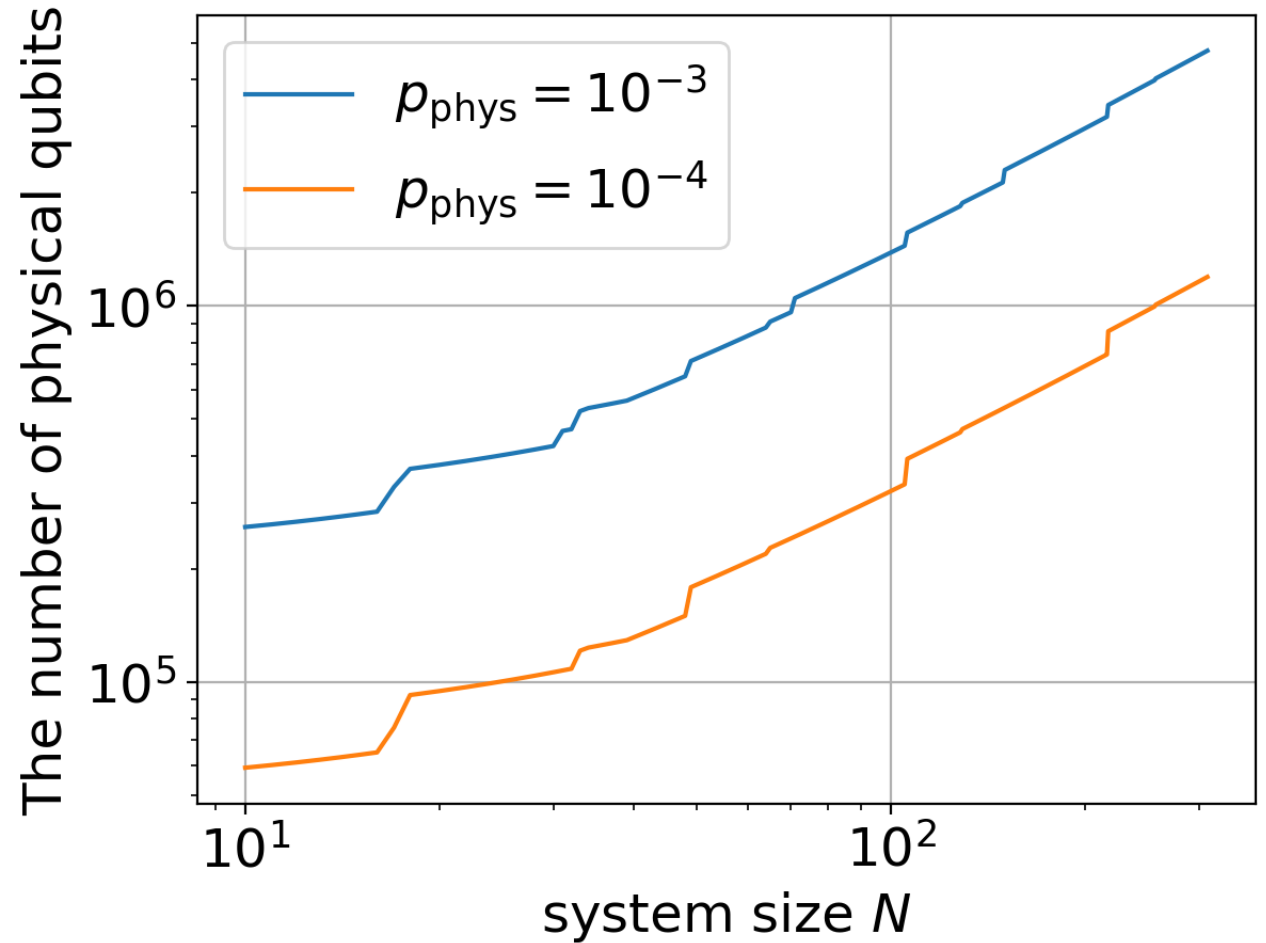
The number of physical qubits

- Parameters

- Precision (additive error) : $\varepsilon = 0.01$
- Evolution time : $t = 4$
- Lattice spacing : $a = 0.2$
- electron mass : $m = 0.1$
- $w = \frac{1}{2a} = 2.5$
- $J = \frac{g^2 a}{2} = 0.1, (g = 1)$
- $\theta_0 = \pi$

Examples ($N = 64$)

Physical error rate	Physical qubits
10^{-3}	9×10^5
10^{-4}	2×10^5



The number of physical qubits for calculating the vacuum persistence amplitude.

Result – resource estimates

Comparison with the previous work (real time evolution e^{-iHt})

A. F. Shaw, et al, Quantum (2020)

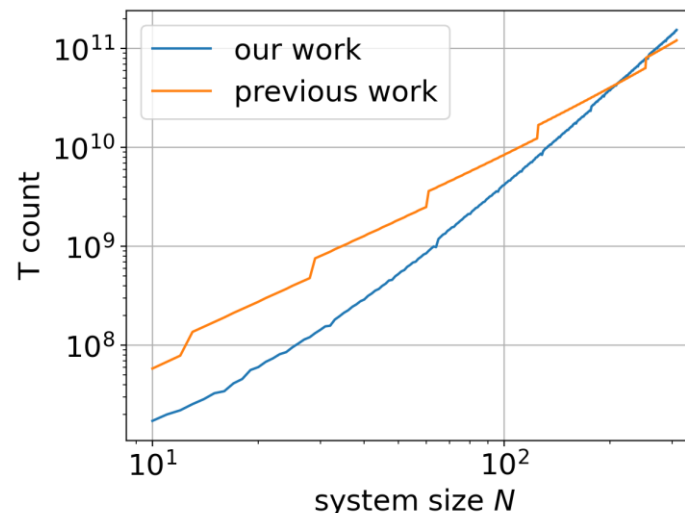
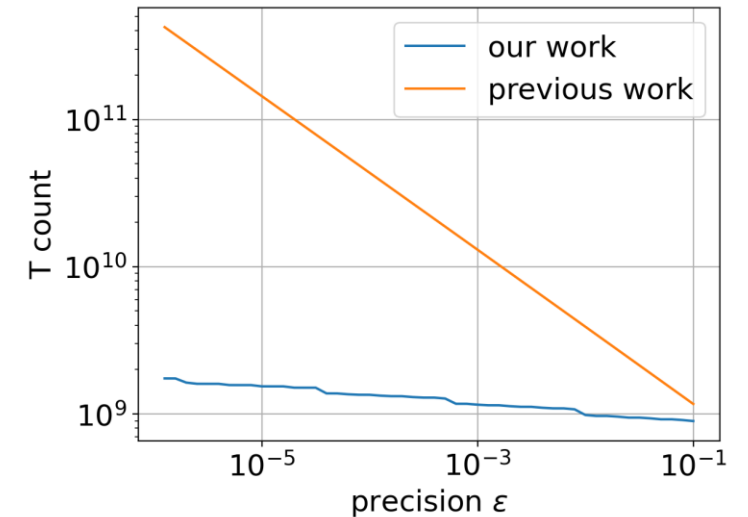
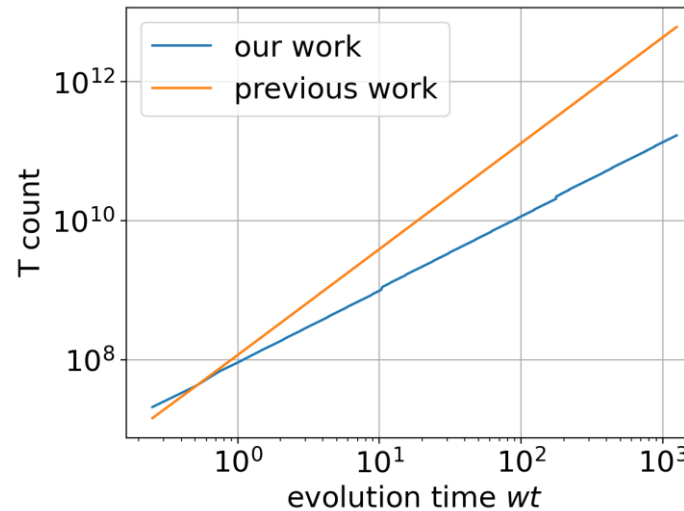
Our work (QSVT) : $\tilde{O}(N^4 t + \log(1/\epsilon))$

The previous work (Trotter)
: $\tilde{O}(N^{2.5} t^{1.5} / \epsilon^{0.5})$

domain	Which is better?
Long-time	Our work
High-precision	Our work
Large-system	Previous work

- Qubit requirement

Our work Previous work
 $N + O(\log N) \ll O(N \log N)$



- Parameters

- Precision (additive error) : $\epsilon = 0.01$
- Evolution time : $t = 4$
- Lattice spacing : $a = 0.2$
- Electron mass : $m = 0.1$
- $w = \frac{1}{2a} = 2.5$
- $J = \frac{g^2 a}{2} = 0.1, (g = 1)$
- $\theta_0 = \pi$

Conclusion

- Comparison to other physical models

Condensed matter physics
(e.g. Hubbard model)

N. Yoshioka, et al, npj Quantum Information 4, 45 (2024)

T count : $\sim 10^8$

<

Schwinger model

$\sim 10^{12}$

\approx

Quantum chemistry
(electronic Hamiltonian)

J. Lee, et al, PRX Quantum (2021)

$\sim 10^{12}$

Summary

- An efficient block-encoding of the Schwinger model Hamiltonian
 - Decompose the Hamiltonian into several parts.
 - Use $O(\log^2 N)$ T gates for P , $O(N)$ T gates for V , with a normalization factor of $O(N^3)$.
- End-to-end complexity for the Schwinger model
 - Estimate the vacuum persistence amplitude.
 - The T gate speed of 1 MHz is minimum requirement.



K. Sakamoto, et. al. Quantum 8, 1474 (2024)

Appendix
