

# Numerical analysis of entanglement entropy using tensor network

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in collaboration with

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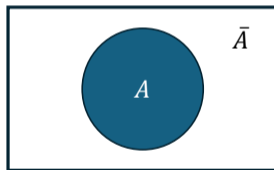
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# Quantum entanglement

- A correlation between two subregions of quantum many-body systems.
- **Entanglement entropy (EE)** is a measure of the degree of quantum entanglement.

$$S_A = -\text{Tr} \rho_A \log \rho_A$$



where  $\rho_A$  is a reduced density matrix of the subregion  $A$  and given by  $\rho_A = \text{tr}_{\bar{A}} \rho$ .

# Entanglement entropy

- Entanglement entropy has many applications in various fields:
  - Particle physics: blackhole entropy, quantum order parameter
  - Quantum information: quantum computer, quantum teleportation
  - Condensed matter physics: phase structure of metallic condensed matter system

# Entanglement entropy

- Entanglement entropy has many applications in various fields:
  - Particle physics: blackhole entropy, **quantum order parameter**
  - Quantum information: quantum computer, quantum teleportation
  - Condensed matter physics: phase structure of metallic condensed matter system

→ We focus on determining the critical temperature of a  $(1+1)$ D lattice model using the EE.

# Entanglement entropy and system detail

The subregion size  $l$  dependence of the EE  $S_A$  tells us the detail of the system.

- Entropic  $c$ -function  $C(l)$ ,  $l$ : length of the subregion  $A$

$$C(l) = \frac{l}{2} \frac{\partial S_A}{\partial l}$$

$C(l)$  monotonically decreases along the RG flow

→ detail of the effective degrees of freedom can be extracted.

- The EE on the quantum critical point is given by

$$S(l) = \frac{c}{3} \log l + k,$$

where  $c$  is a central charge and  $k$  is a constant.

→ central charge can be extracted.

# Numerical analysis of entanglement entropy

- Monte Carlo method
  - calculates the entropic  $c$ -function.
    - e.g. 4D SU(3) gauge theory [Itou-Nagata-Nakagawa-Nakamura-Zakharov, 2015]
  - based on the definition of the EE on lattice
    - [Aoki-Iritani-Nozaki-Numasawa-Shiba-Tasaki, 2015],
  - has **sign problem**.

# Numerical analysis of entanglement entropy

- Tensor network method (In this talk, we focus on those of Lagrangian formalism)
  - has no sign problem.
  - directly computes the reduced density matrix and the EE.  
e.g. (1+1)D  $O(3)$  non-linear sigma model [Kuramashi-Luo, 2023]
  - is limited to the case of half-space subregion  $A$ .

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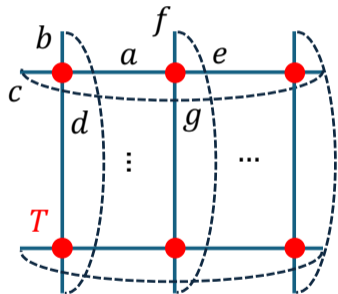
We compute the subregion size dependence of the EE with our new method  
[Hayazaki-Kadoh-Takeda-GT, work in progress].



# Tensor network (Lagrangian formalism)

- Partition function can be directly computed.

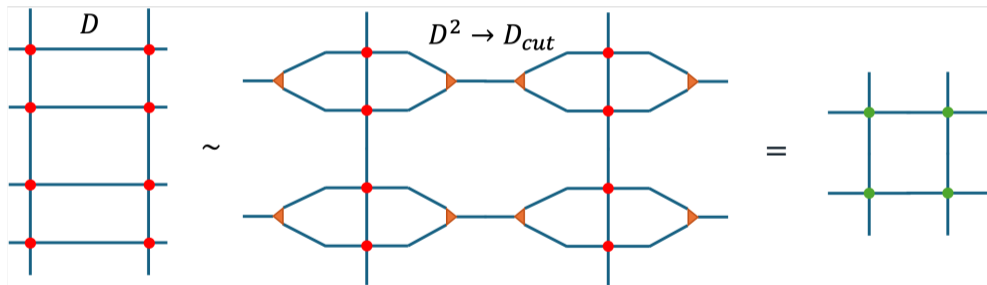
$$Z = \int \mathcal{D}\phi e^{-S[\phi]} = \sum_{\dots, a, b, c, d, e, f, g, \dots} \dots T_{abcd} T_{efag} \dots =$$



- We need some "coarse-graining" to reduce the computational cost.

# Tensor Renormalization group (TRG)

- Recursively approximates multiple tensors as one tensor.  
e.g. Higher-order TRG (HOTRG) algorithm

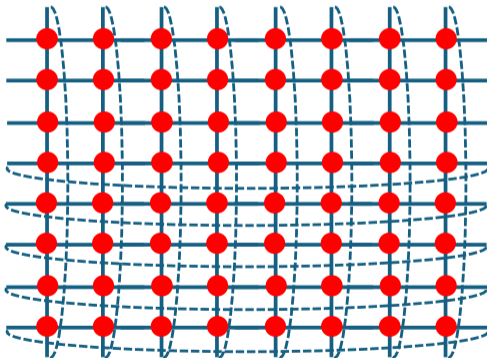


- Various TRG algorithms are proposed:  
A-TRG [Adachi-Okubo-Todo, 2019, 2019], Triad-TRG [Kadoh-Nakayama, 2019], etc.



# Tensor network computation of reduced density matrix

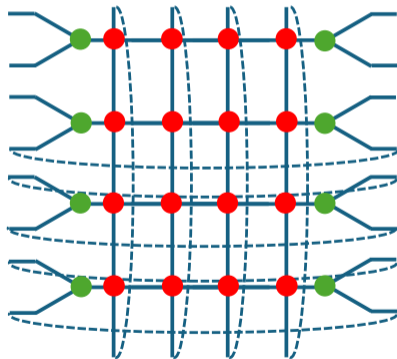
Example: total spatial size is 8 and subregion size 3.



Tensor network representation of the reduced density matrix before coarse-graining.

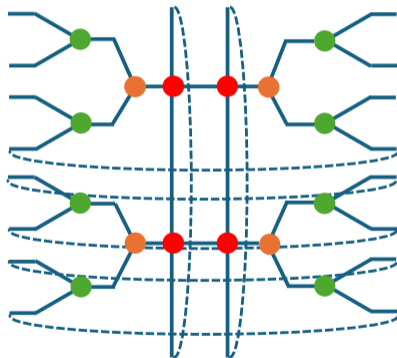
# Tensor network computation of reduced density matrix

After one HOTRG coarse-graining procedure:



# Tensor network computation of reduced density matrix

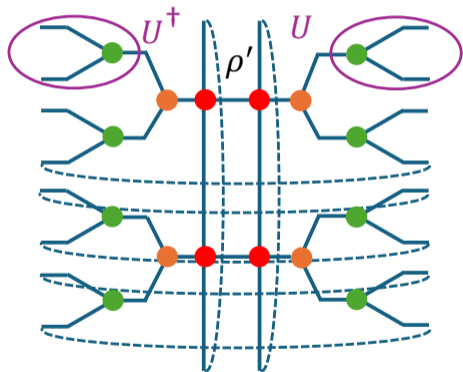
After two HOTRG coarse-graining procedures:



At this stage, we can simplify this network further!

# Tensor network computation of reduced density matrix

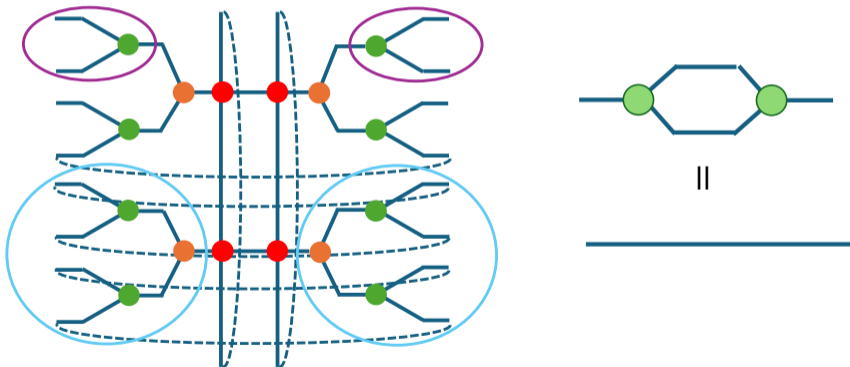
Tensors  $U$  and  $U^\dagger$  do not contribute to the entanglement entropy.



$$S_A = -\text{tr} \rho_A \log \rho_A = -\text{tr} U^\dagger \rho'_A U \log (U^\dagger \rho'_A U) = -\text{tr} \rho'_A \log \rho'_A$$

# Tensor network computation of reduced density matrix

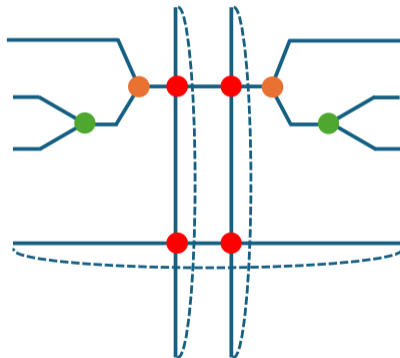
Some isometry tensors can be contracted and become an identity matrix.





# Tensor network computation of reduced density matrix

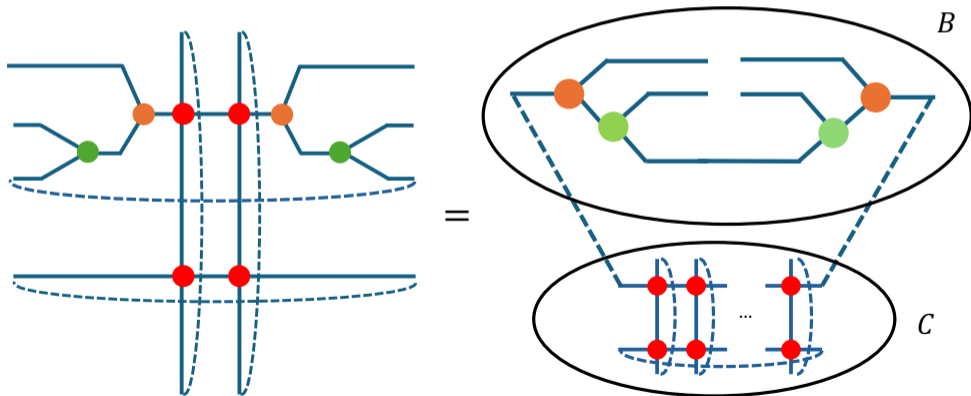
Finally, we obtain the simplified tensor network of the reduced density matrix below:



We established the algorithm to obtain this final result directly.

# Our algorithm

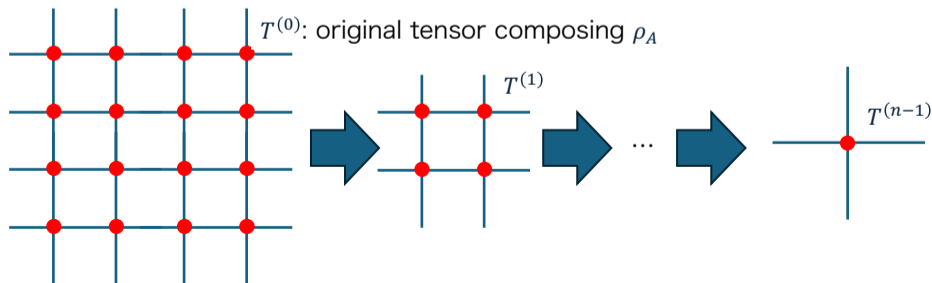
The simplified tensor network of the reduced density matrix consists of two parts: core matrix  $C$  and boundary factor  $B$ .



# Our algorithm

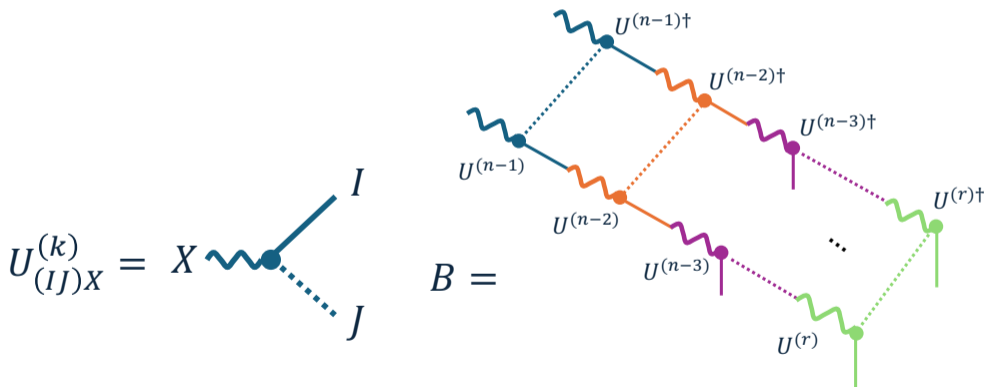
In the following, we set the total spatial size  $L_x = 2^n$ , temporal size  $L_t = \alpha \cdot 2^n$ .  
The core matrix  $C$  consists of coarse-grained tensor  $T^{(n-1)}$ .

$$C_{(x_1 x_2)(x'_1 x'_2)} = \begin{array}{c} T^{(n-1)} \\ \begin{array}{c} x_1 \quad \bullet \quad \bullet \quad \bullet \quad x'_1 \\ x_2 \quad \bullet \quad \bullet \quad \bullet \quad x'_2 \end{array} \\ \underbrace{\hspace{10em}}_{2\alpha} \end{array}$$



# Our algorithm

The boundary factor  $B$  consists of isometry tensors  $U^{(i)}$  and  $U^{(i)\dagger}$  obtained in the coarse-graining procedure of tensor  $T^{(n-1)}$



The contraction of isometry tensors depends on the subregion size  $l$ .

# Numerical Analysis: (1+1)D XY model

- Partition function and action:

$$Z = \int \prod_{x=0}^{L_x} \prod_{t=0}^{L_t} \frac{d\theta_{x,t}}{2\pi} e^{-S}$$

$$S = -\beta \sum_{x,t} \cos(\theta_{x,t+1} - \theta_{x,t}) - \beta \sum_{x,t} \cos(\theta_{x+1,t} - \theta_{x,t})$$

$\beta$ : inverse temperature

Spatial lattice size: 1024, temporal lattice size:  $2^8 \times 1024$

- XY model exhibits the topological BKT phase transition at  $T = T_{\text{BKT}}$ , and  $0 < T < T_{\text{BKT}}$  is the critical line. ( $T_{\text{BKT}} = 0.892943(2)$  [Ueda-Oshikawa, 2021])

# Tensor network for XY model

Partition function  $Z$ :

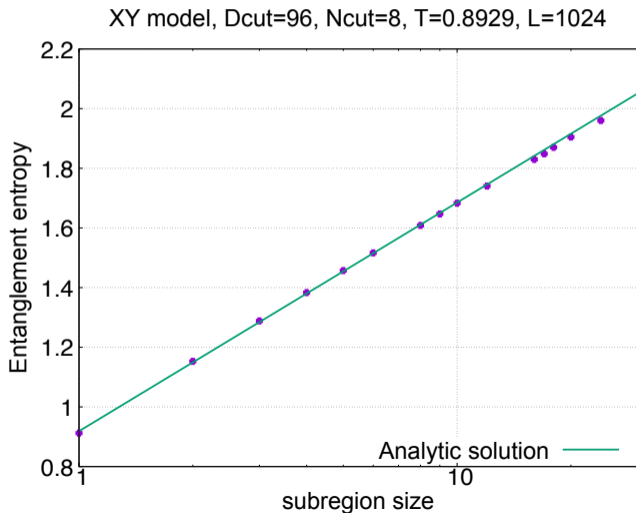
$$Z = \int \prod_{x=0}^{L_x} \prod_{t=0}^{L_t} \frac{d\theta_{x,t}}{2\pi} e^{-S} = \prod_{\text{lattice}} T_{xx'yy'}$$

$$T_{xx'yy'} \equiv \sqrt{e^{(y+y')\mu}} \delta_{x'+y'-x-y} \sqrt{I_{y'}(\beta)} \sqrt{I_y(\beta)} \sqrt{I_{x'}(\beta)} \sqrt{I_x(\beta)}$$

$I_x(\beta)$ : modified Bessel function of the first kind, where  $x$  takes from  $-\infty$  to  $\infty$ .

→ We regularize  $I_x(\beta)$  by introducing the cutoff  $N_{\text{cut}}$ :  $-N_{\text{cut}} \leq x \leq N_{\text{cut}}$

# Result - subregion size dependence of EE and central charge



- subregion size  $l$ :  
 $l = 2^p + 2^q$  ( $q < p$ )
- Analytic solution of EE of finite size subregion

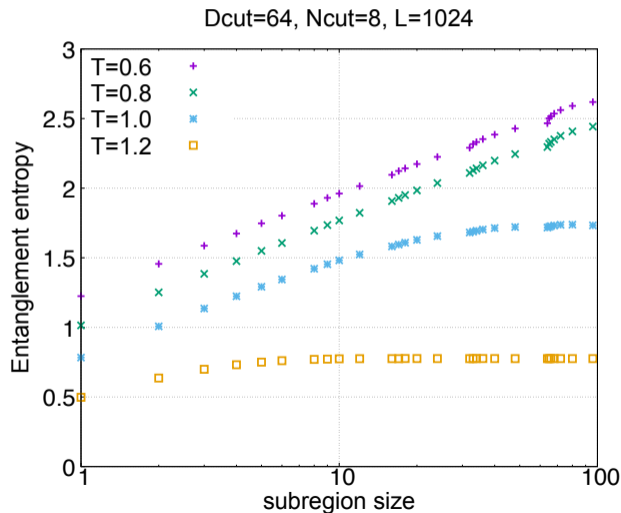
$$S(l, L) = \frac{c}{3} \log \left( \sin \left( \frac{\pi l}{L} \right) \right) + k$$

- Central charge  $c$  by fitting the result to the analytic solution

$$c = 0.998(5)$$

→ agrees with known result  $c = 1$ .

## Result - temperature dependence of EE



- On the critical line

$$T = 0.6, 0.8 < T_{BKT}$$

$$S(l, L) = \frac{c}{3} \log \left( \sin \left( \frac{\pi l}{L} \right) \right) + k$$
$$\sim \frac{c}{3} \log l + \text{const.}$$

- Non-critical  $T = 1.0, 1.2 > T_{BKT}$ :  
 $l$  dependence for small  $l$   
 $\therefore$  finite correlation length.
- We may determine the transition temperature using the EE.



### Summary of this talk:

- We studied the subregion size dependence of the entanglement entropy in the 1+1D XY model.
- We determined the central charge on the critical line  $T < T_{\text{BKT}}$  using the subregion size dependence of the EE.
- Difference in the behavior of the EE implies that we can determine the transition temperature using the EE.

# Conclusion and discussion

Future direction:

- Compute entanglement entropy of a larger subregion size.
- Determine transition temperature.

Method:

- More efficient TRG algorithm  
e.g. HOSRG [Z. Y. Xie, et al., 2012]
- Parallelization of algorithm  
e.g. Parallelized HOTRG [Yamashita-Sakurai, 2021]

## Backup - Determining the fitting range

- 1 Compute the EE  $S(l')$ ,  $S(l)$  and obtain the central charge  $c(l)$  where  $l = 2^p$ ,  $l' = 2^{p+1}$  or  $l = 2^p + 2^q$ ,  $l' = 2^{p+1} + 2^q$ .

$$c(l) = 3 \frac{S_A(l') - S_A(l)}{\log \sin \frac{l'\pi}{1024} - \log \sin \frac{l\pi}{1024}}$$

$$c(1) = 1.041435054$$

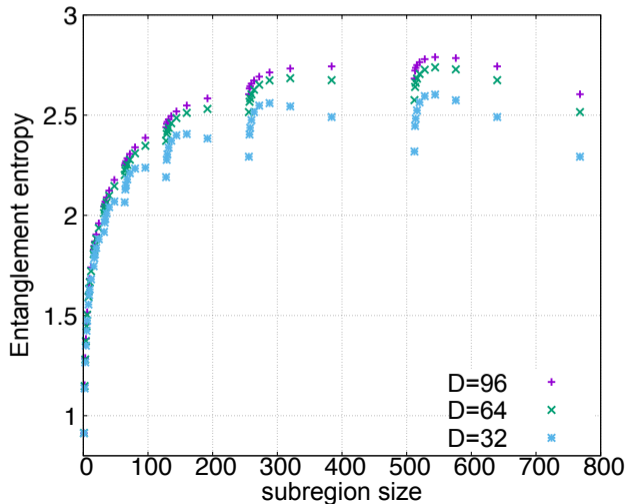
$$c(2) = 0.997233754, \dots$$

- 2 Determine the fitting range as 1 to  $x$ , where  $x$  is the smallest integer that satisfies

$$|c(x) - 1| \geq |c(1) - 1| = 0.041435054.$$

# Backup - Dcut dependence of the EE

XY model, Ncut=8, aspect ratio  $2^8$ ,  $T=0.8929$ ,  $L=1024$



## Backup - Boundary factor

The boundary factor  $B$  is composed of isometries  $U^{(n-2)}, U^{(n-3)}, \dots, U^{(r)}$ .

The integer  $r$  is the largest one that satisfies  $a_k \neq b_k$ , where

$$l = \sum_{k=0}^{n-1} a_k 2^k \quad (a_k = 0, 1),$$
$$l - 1 = \sum_{k=0}^{n-1} b_k 2^k \quad (b_k = 0, 1).$$

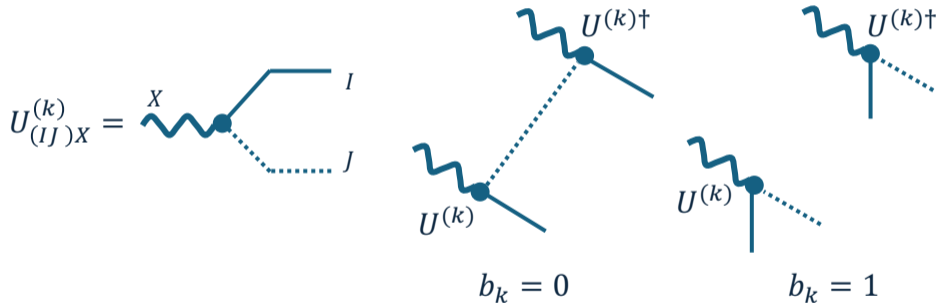
For example, letting  $L = 2^4$  and  $l = 5$ , we have

$$l = 0 \cdot 2^3 + 1 \cdot 2^2 + 0 \cdot 2^1 + 1 \cdot 2^0,$$
$$l - 1 = 0 \cdot 2^3 + 1 \cdot 2^2 + 0 \cdot 2^1 + 0 \cdot 2^0,$$

and  $r = 0$ .

## Backup - Boundary factor

$b_k$  determines the form of contraction of isometry  $U^{(k)}$  and  $U^{\dagger(k)}$ .

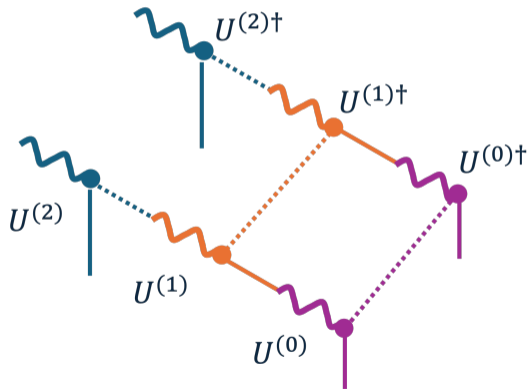


The index of  $U^{(k)}$  represented by a wavy line is contracted with the index of  $U^{(k+1)}$  represented by a solid line or a dotted line.

## Backup - Boundary factor

Example: Total spacial size 16, and subregion size 5.

$\rightarrow b_2 = 1, b_1 = 0, b_0 = 0$  and  $r = 0$ .



The indices represented by a wavy line are contracted with core matrix.